The Quantum Mechanical Treatment of Viscosity by Use of the Rigid Elastic Sphere Model. II. The Sutherland Constant.

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(Received April 4, 1944)

Introduction. In part one of the present study in which the viscosity  $\eta$  was treated theoretically from the standpoint of quantum mechanics, it was shown that its temperature variation was reproduced more correctly over the wide ranges than from the classical standpoint, and especially so if the term originated by Sutherland was introduced, when the elastic rigid sphere model was used. These consequences led us to consider that the Sutherland constant C, determined empirically by the classical formula of viscosity:

$$\eta \propto \frac{\sqrt{T^3}}{Q_{\eta 0}(1+C/T)}$$
, (1)

where  $Q_{\eta_0}$  denoted the viscosity close-section, included not only the gen-

erally accepted meaning of the degree of the molecular attraction,  $^{(2)}$  but also the quantum mechanical effect. From such reasons it becomes evident that some of the toichiometrical studies dealing with the constant would be worthy of reinspection by use of the new constant  $C_0$  which correlates to the degree of the intermolecular potential only. It is the object of the present study to carry it out, and to examine the above conclusion.

Determination of  $C_0$ . The constant  $C_0$  can be obtained, as in the case of neon, by such procedure that the temperature variation of  $\eta$  fits well with the theoretical formula in the previous paper. However, as it is too laborous to carry out the procedure for many substances, the following simplified method was adopted conventionally.

It was assumed that the quantal viscosity cross-section did not depend upon the velocity of molecules but depend only on the temperature. Thus the Sutherland constant  $C_0$  having quantal meaning was given also by (1), on the condition that  $Q_{\eta_0}$  and C must be replaced by  $Q_{\eta}$  and  $C_0$ , respectively, in the formula. So the  $C_0$  can be determined by the relation:

$$T = H \frac{\sqrt{T^3}}{\eta_a} \dot{-} C_0 \tag{2}$$

graphically, where H denotes the proportional constant.

Sutherland Constant and Boiling-point. It is known that the Sutherland constant of various substances has a close relation with their boiling-points, and consequently also with their critical points. These relations are expressed by the following simple rules:

Rankine's rule . . . . 
$$C = T_k/1.12$$
,  
Vogel's rule . . . .  $C = 1.47 \times T_S$ ,

where  $T_K$  and  $T_S$  denote the critical and the boiling points expressed in absolute temperature, respectively.

As shown in table 1, these rules reproduce the experimental values fairly well in most of the substances. The theoretical derivation of them has not been succeeded, though they are presumed plausible from the qualitative point of view, showing the correlations of the intermolecular force between the concerned two quantities. To the exceptions belong such gases as acetylene, hydrogen, helium, hydrogen chloride and water. (5) However, if we examine the table carefully, a tendency of systematic deviation is recognized; i.e., some substances which have smaller molecular weight give too low values as expected from the rules. Now, the discrepancy found in the low molecular gases may be removed, if we substitute the constant C for the new one  $C_0$  calculated as above, considering the

and

<sup>(1)</sup> See the previous paper.

<sup>(2)</sup> Kennard, "Kinetic Theory of Gases", New York (1938) 154.

<sup>(3)</sup> Rankine, Proc. Roy. Soc., A84(1910), 181.

<sup>(4)</sup> Vogel, Ann. der Physik, 43 (1914), 1235.

<sup>(5)</sup> Most of the figures are indebted to the table given by H. Adzumi ("Kwagaku-Jikkengaku", 2(1942), 40).

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meaning of the rules. The idea was tested as shown in the third column of table 1.

Table 1. Test of the Rankine's and Vogel's Rules.

		Sutherland Constand	Calculated by	$T_K$	$1.47 \times T_S$
	C	lassical Method Q	uantal Method	1.12	202000
He	{ 80 70	(Schultz) (Vogel)	32 (20–100°C)	4.6	6.2
Ne	$\left\{\begin{array}{c} 61 \\ 56 \end{array}\right.$	(Trautz-Binkele) (Rankine)	42 (15-100°C)	40	40
A	$\left\{\begin{array}{c}142\\147\end{array}\right.$	(Rankine) (Trautz-Binkele)	136 (20- 80°C)	135	124
Kr	188	(Rankine)		187	178
$\mathbf{H}_2$	$\left\{\begin{array}{c} 83\\83\\79\end{array}\right.$	(Schmidt) (Trautz-Stauf) (Adzumi)	31 (20-100°C)	30	30
$\mathbf{D_2}$	79*	k	32 (-50- 23°C)		24
$N_2$	$\left\{ \begin{array}{c} 103 \\ 118 \end{array} \right.$	(Trautz-Baumann) (Smith)		112	114
$O_2$	$\left\{ egin{array}{l} 138 \\ 125 \end{array} \right.$	(Vogel) (Trautz-Melster)		137	132
$\mathrm{CH_4}$	$\left\{ \begin{array}{c} 170 \\ 198 \end{array} \right.$	(Adzumi) (Rankine)	147 (40–100°C) 175 ( 0–100°C)	170	160
$C_2H_6$	$\left\{ egin{array}{l} 289 \ 287 \end{array}  ight.$	(Adzumi) (Titani)	267 (20–100°C) 260 (20–100°C)	275	<b>2</b> 72
$C_2H_4$	259	(Titani)	247 (20-120°C)	253	249
$C_3H_8$	318	(Titani)		330	335
· HCl	$\left\{\begin{array}{c} 317\\369\end{array}\right.$	(Harle) (Trautz-Narath)		289	279
$H_2O$	$\left\{\begin{array}{c} 961 \\ 650 \end{array}\right.$	(Braune-Linke) · (Schuhgajew)		576	548
$C_2H_2$	$\left\{ egin{array}{l} 196 \\ 220 \end{array}  ight.$	(Titani) (Adzumi)	· <u></u>	276	278
СО	$\left\{ \begin{array}{c} 118 \\ 111 \end{array} \right.$	(Smith) (Trautz-Bauman)		120	122
$CO_2$	274	(Smith)		271	286

<sup>\*</sup> Calculated from the data of Van Cleave and O. Masse (J. Can. Res., 12 (1935), 57).

According to the above procedure,  $C_0$ s were calculated for several gases, the molecular weights of which were comparatively small. As shown in table 1, the coincidence of the values with the experimental ones became better in most cases, (6) except such gases as acetylene and water, these exception being probably on account of the same reason as occurred in Trouton's rule. However, the improvement is slight for helium;  $C_0$  is still too large. Such circumstances may be due to the insufficiency of the correction term of Sutherland type, because, if we use the values at

<sup>(6)</sup> The formula is given by equation (1) times  $1/(1+C_0/T)$ .

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low temperatures, the discrepancy can be improved; i.e.,  $C_0=4$  from 273-89°K, and  $C_0$  tends to zero from 89° downwards.<sup>(7)</sup>

Generally speaking, the coincidences seem better if the constants 1.12 and 1.47 were selected slightly larger and smaller, respectively, because these were determined in connection with  $C_0$ , and not with  $C_0$ .

The Sutherland Constant and Electric Moment. As the Sutherland constant is a measure of potential energy between molecules, it can be derived theoretically for simple gases. For nonpolar gases, Debye<sup>(8)</sup> has obtained the following formula:

The Sutherland constant = 
$$\frac{217}{205} \cdot \frac{7}{20\pi} \cdot \frac{P_0 \tau^2}{R r_0^8}, \qquad (3)$$

where  $P_0$  is the molecular polarisation,  $\tau$  the mean electric moment and R the gas constant. For some quadrapole molecules the equation was tested by himself, comparing the observed  $\tau$  with those calculated from other values. As shown in table 2, it was found that there lies some disagreement in lighter gases such as in hydrogen and helium. However, as the Sutherland constant which was used for the calculation was C and not  $C_0$  if the notation of this paper was used, considerable disagreement has been observed for the lighter gases, because C has the meaning as pointed above by the author. This idea may be ascertained if the calculation is carried out, using  $C_0$  instead, and the other constants being reserved unvaried. As shown in table 2, the calculated electric moment for He,  $H_2$  and Kr (the fourth column) coincides with the experimental values (the second column), thus the result as expected being obtained.

It becomes thus evident that the Sutherland constant which is determined by the quantum-mechanical formula can be regarded as a reasonable measure of the intermolecular force. This conclusion becomes, in turn, a support to the argument in the introduction of the previous paper; i.e., the necessity of the application of the quantum mechanics in the transport phenomena.

	τ from the	$\tau$ from eq. (1) and from		
Gase	critical data	C obtained classically	$C_0$ obtained quantally	
He	2.84	4.33(75)**	2.83(32)*	
$\mathbf{Kr}$	11.2	11.4 (162)	10.5 (136)	
$\mathbf{X}\mathbf{e}$	18.9	20.6 (252)		
$\mathbf{H_2}$	3.20	5.60(182)	3.4 (31)	
$N_2$	13.3	13.4 (113)		
$O_2$	11.2	11.7 (137)		

Table 2. Electric Moment and Sutherland Constant.

<sup>\*</sup> Considering the temperature range used for comparison, 32 was adopted.

<sup>\*\*</sup> Numbers in parenthese are the Sutherland constants used.

<sup>(7)</sup> As it may be expected, the determination of C became easier in this method, because the experimental values lie on a straight line over wider temperature ranges.

<sup>(8)</sup> P. Debye, Physik. Zeit., 21 (1921), 179.

## Summary.

- 1) According to the conclusion of the previous paper, the "Sutherland constant" in the viscosity coefficient for various substances has been calculated quantum-mechanically by a simple procedure.
- 2) The Sutherland constant  $C_0$  thus calculated has been found to obey Rankin's and Vogel's rules even in the substances of low molecular weight.
- 3) The disagreement in the Debye's formula between electric moment and Sutherland constant has been shown to disappear if  $C_0$  was used instead of the classically determined Sutherland constant.

The author wishes to express his sincere thanks to Dr. T. Arimori and to Dr. Y. Tanaka of our laboratory for their encouragement throughout this research.

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